

Q1. Write down the operators  $U_R(\delta\phi_x)$  and  $U_R(\delta\phi_y)$  for infinitesimal rotations about the Cartesian axes X and Y through infinitesimal rotation angles and find their commutator.

→2 marks

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A1:  $U_R(\delta\phi_x) = 1 - \frac{i}{\hbar} \delta\phi_x J_x$

$$U_R(\delta\phi_y) = 1 - \frac{i}{\hbar} \delta\phi_y J_y$$

The operators do not commute since for this operator, rotation through X axis and then through Y axis is not the same when reversed.

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Q2. In the Pauli-Lenz quantum vector for the hydrogen atom given by  $\vec{A}_{QM} = \frac{1}{2\mu} [\vec{p} \times \vec{L} - \vec{L} \times \vec{p}] - \kappa \hat{r}$ , express  $\kappa$  in terms of the physical properties of the electron and state the dimensions of  $\kappa$ .

→2 marks

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A2.  $\kappa$  is the parameter which gives the strength of the centre field potential. The dimensions of  $\kappa$  is  $[\vec{v} \times \vec{H}] = LT^{-1} \times L^2T^{-1} = L^3T^{-2}$

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Q3. Recall that  $[J_+, J_-]_- = 2\hbar J_z$  and  $J_+ |jm\rangle = \lambda_m \hbar |j, m+1\rangle$  using these relations find  $|\lambda_{m-1}|^2 - |\lambda_m|^2$  in terms of the eigenvalue  $\hbar m$  of the operator  $J_z$ .

→3 marks

A3.

$$[J_+, J_-]_- = 2\hbar J_z$$

$$\langle jm | J_+ J_- - J_- J_+ | jm \rangle = 2\hbar \langle jm | J_z | jm \rangle$$

$$\langle jm | J_+ J_- | jm \rangle - \langle jm | J_- J_+ | jm \rangle = 2\hbar \langle jm | J_z | jm \rangle$$

$$\lambda_{m-1}^* \hbar \langle jm | J_+ | jm-1 \rangle - \lambda_m \hbar \langle jm | J_- | jm+1 \rangle = 2\hbar \langle jm | J_z | jm \rangle$$

$$\lambda_{m-1}^* \hbar \lambda_{m-1} \hbar \langle jm || jm \rangle - \lambda_m \hbar \lambda_m^* \hbar \langle jm || jm \rangle = 2\hbar m \hbar \langle jm || jm \rangle$$

$$|\lambda_{m-1}|^2 - |\lambda_m|^2 = 2m$$

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Q4. When can a vector  $\vec{A}$  be a vector operator  $\vec{A}_{op}$ ? Explain how the condition for  $\vec{A}$  to be the vector operator  $\vec{A}_{op}$  determines how the angular momentum of a system can be defined.

→2 marks

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A4. A vector  $\vec{A}$  can be a vector operator  $\vec{A}_{op}$  if it transforms under rotation as per

$$U_R^\dagger A_i U_R = \sum_{j=1}^3 R_{ij} A_j \text{ or equivalently it should follow the commutation relations } [A_i, J_j] = i\hbar \epsilon_{ijk} A_k .$$

where  $\vec{J} = \sum_{j=1}^3 J_j \hat{e}_j$  is the generator of rotations. It is called as angular momentum.

The transformation of vector A to be an  $\vec{A}_{op}$  determines the definition of angular momentum to be the generator of rotations such that the above relations hold good. (If not either the vector A cannot be a vector operator or the angular momentum operator must be defined differently)

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Q5. What are the principal objectives of the Foldy-Wouthuysen transformations of the Dirac Hamiltonian? → 1 mark

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A5. The Principle objective is to decouple the Dirac equation into *two* two-component equations: one reduces to the Pauli description in the non-relativistic limit; the other describes the negative energy states.

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Q6. Prove that: 
$$\frac{\left( c\vec{\alpha} \cdot \left( \vec{p} - \frac{e}{c} \vec{A} \right) \right)^2}{2mc^2} = \frac{1}{2m} \left\{ \vec{\pi}^2 + i\vec{\sigma} \cdot \vec{\pi} \times \vec{\pi} \right\} 1_{4 \times 4}$$
 →3 marks

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A6. 
$$\frac{\left( c\vec{\alpha} \cdot \left( \vec{p} - \frac{e}{c} \vec{A} \right) \right)^2}{2mc^2} = \frac{1}{2m} \left\{ \vec{\alpha} \cdot \left( \vec{p} - \frac{e}{c} \vec{A} \right) \right\} \left\{ \vec{\alpha} \cdot \left( \vec{p} - \frac{e}{c} \vec{A} \right) \right\}$$

$$\begin{aligned} \frac{\theta^2}{2mc^2} &= \frac{1}{2m} \left\{ \begin{pmatrix} \vec{0} & \vec{\sigma} \\ \vec{\sigma} & \vec{0} \end{pmatrix} \cdot \vec{\pi} \right\} \left\{ \begin{pmatrix} \vec{0} & \vec{\sigma} \\ \vec{\sigma} & \vec{0} \end{pmatrix} \cdot \vec{\pi} \right\} \\ &= \frac{1}{2m} \begin{pmatrix} \vec{0} & \vec{\sigma} \cdot \vec{\pi} \\ \vec{\sigma} \cdot \vec{\pi} & \vec{0} \end{pmatrix} \begin{pmatrix} \vec{0} & \vec{\sigma} \cdot \vec{\pi} \\ \vec{\sigma} \cdot \vec{\pi} & \vec{0} \end{pmatrix} \end{aligned}$$

Where  $\theta = c\vec{\alpha} \cdot \left( \vec{p} - \frac{e}{c} \vec{A} \right)$  and  $\pi = \left( \vec{p} - \frac{e}{c} \vec{A} \right)$

$$\begin{aligned} \frac{\theta^2}{2mc^2} &= \frac{1}{2m} \begin{pmatrix} \vec{\sigma} \cdot \vec{\pi} \vec{\sigma} \cdot \vec{\pi} & 0 \\ 0 & \vec{\sigma} \cdot \vec{\pi} \vec{\sigma} \cdot \vec{\pi} \end{pmatrix}; \text{ Using Pauli's identity} \\ &= \frac{1}{2m} \vec{\sigma} \cdot \vec{\pi} \vec{\sigma} \cdot \vec{\pi} 1_{4 \times 4} \end{aligned}$$


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$$\frac{\theta^2}{2mc^2} = \frac{1}{2m} \{ \vec{\pi}^2 + i\vec{\sigma} \cdot \vec{\pi} \times \vec{\pi} \} 1_{4 \times 4}$$


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Q7. Determine if the exchange integral for a pair of electrons can be non-zero if their spins are anti-parallel. →2 marks

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The Exchange integral,

$$K = \iint dq_1 dq_2 u_{\alpha_i}^*(q_2) u_{\alpha_j}^*(q_1) \frac{1}{r_{12}} u_{\alpha_i}(q_1) u_{\alpha_j}(q_2)$$

$$\text{Now, } \int dq \text{ ————— } \equiv \iiint dV \text{ ————— } \sum_{\zeta} \text{ — }$$

$$K = \iint dV_1 dV_2 u_{\alpha_i}^*(\vec{r}_2) u_{\alpha_j}^*(\vec{r}_1) \frac{1}{r_{12}} u_{\alpha_i}(\vec{r}_1) u_{\alpha_j}(\vec{r}_2) \times \sum_{\zeta_1} \sum_{\zeta_2} \langle \zeta_2 | m_{s_i} \rangle^* \langle \zeta_1 | m_{s_j} \rangle^* \langle \zeta_1 | m_{s_i} \rangle \langle \zeta_2 | m_{s_j} \rangle$$

$$K = \iint \text{ ————— } \times \sum_{\zeta_1} \sum_{\zeta_2} \langle m_{s_i} | \zeta_2 \rangle \langle m_{s_j} | \zeta_1 \rangle \langle \zeta_1 | m_{s_i} \rangle \langle \zeta_2 | m_{s_j} \rangle$$

$$K = \iint \text{ ————— } \times \langle m_{s_i} | m_{s_i} \rangle \times \langle m_{s_j} | m_{s_j} \rangle$$

$$K = \iint \text{ ————— } \times \delta_{m_{s_j} m_{s_i}}$$

⇒ Exchange integral exists only if spins are antiparallel.

