

Boundary Layer Theory (Lectures 37 to 40)

Q1. Choose the correct answer

- (i) The boundary layer thickness for flow over a flat plate
- (a) increases with an increase in the free stream velocity
 - (b) decreases with an increase in the free stream velocity
 - (c) increases with an increase in the kinematic viscosity
 - (d) decreases with an increase in the kinematic viscosity

[Ans. (b) and (c)]

- (ii) If x is the distance measured from the leading edge of a flat plate, the wall shear stress for laminar boundary layer varies as

- (a) x
- (b) $x^{1/2}$
- (c) $x^{-1/2}$
- (d) $x^{-4/5}$

[Ans.(b)]

- (iii) The growth of the turbulent boundary layer thickness as compared to the laminar boundary layer takes place

- (a) at a slower rate
- (b) at a faster rate
- (c) at the same rate
- (d) unpredictable

[Ans.(b)]

- (iv) Separation in flow past a solid object is caused by

- (a) a favourable (negative) pressure gradient
- (b) an adverse (positive) pressure gradient
- (c) the boundary layer thickness reducing to zero
- (d) a reduction of pressure to vapour pressure

[Ans.(b)]

Q2.

The velocity profile over a flat plate of length L is expressed in terms of a similarity

variable η as $\frac{u}{U_\infty} = \frac{dF(\eta)}{d\eta}$, where $\eta = y\sqrt{\frac{U_\infty}{\nu x}}$. Numerical solution for the variable

reveals that $\left. \frac{d^2F}{d\eta^2} \right|_{\eta=0} = 0.3$. Determine the drag coefficient as a function of Reynolds

number, Re_L .

Solution

The wall shear stress is found to be

$$\tau_w = \mu \left. \frac{\partial u}{\partial y} \right|_{y=0}$$

$$\begin{aligned}
&= \mu \left. \frac{\partial u}{\partial \eta} \right|_{\eta=0} \frac{\partial \eta}{\partial y} \\
&= \mu U_{\infty} \left. \frac{d^2 F}{d\eta^2} \right|_{\eta=0} \sqrt{\frac{U_{\infty}}{\nu x}} \\
&= 0.3\mu U_{\infty} \sqrt{\frac{U_{\infty}}{\nu x}}
\end{aligned}$$

The drag force acting on one side of the plate is

$F_D = \int_0^L \tau_w w dx$ (w is the width of the plate in a direction perpendicular to the plate)

$$\begin{aligned}
&= \int_0^L 0.3\mu U_{\infty} \sqrt{\frac{U_{\infty}}{\nu x}} w dx \\
&= 0.6U_{\infty} w \mu \sqrt{\text{Re}_L}
\end{aligned}$$

The drag coefficient is then

$$C_{fL} = \frac{F_D}{\frac{1}{2}\rho U_{\infty}^2 w L} = \frac{0.6U_{\infty} w \mu \sqrt{\text{Re}_L}}{\frac{1}{2}\rho U_{\infty}^2 w L} = \frac{1.2}{\sqrt{\text{Re}_L}}$$

Q3.

The velocity profile within boundary layer for steady, two-dimensional, constant density, laminar flow over a flat plate is a polynomial of order 4 and is given as:

$$\frac{u}{U_{\infty}} = C_0 + C_1 \frac{y}{\delta} + C_2 \left(\frac{y}{\delta}\right)^2 + C_3 \left(\frac{y}{\delta}\right)^3 + C_4 \left(\frac{y}{\delta}\right)^4$$

Using suitable boundary conditions, evaluate the constants C_0 , C_1 , C_2 , C_3 and C_4 .

Solution

The boundary conditions are

(i) at $y = 0$, $u = 0$

(ii) at $y = \delta$, $u = U_{\infty}$

(iii) at $y = \delta$, $\frac{\partial u}{\partial y} = 0$

(iv) at $y = 0$, $\frac{\partial^2 u}{\partial y^2} = 0$

(v) at $y = \delta$, $\frac{\partial^2 u}{\partial y^2} = 0$

Applying $u = 0$ at $y = 0$, one can get

$$0 = C_0$$

Applying $u = U_{\infty}$ at $y = \delta$ gives

$$1 = C_1 + C_2 + C_3 + C_4 \quad (1)$$

Differentiating the velocity profile (note that both u and δ are functions of x only) with respect to y , we obtain

$$\frac{1}{U_\infty} \frac{\partial u}{\partial y} = \frac{C_1}{\delta} + \frac{2C_2}{\delta} \frac{y}{\delta} + \frac{3C_3}{\delta} \left(\frac{y}{\delta}\right)^2 + \frac{4C_4}{\delta} \left(\frac{y}{\delta}\right)^3$$

Applying (iii) boundary condition i.e., $\frac{\partial u}{\partial y} = 0$ at $y = \delta$, we get

$$0 = C_1 + 2C_2 + 3C_3 + 4C_4 \quad (2)$$

A second differentiation of the velocity profile gives

$$\frac{1}{U_\infty} \frac{\partial^2 u}{\partial y^2} = \frac{2C_2}{\delta^2} + \frac{6C_3}{\delta^2} \frac{y}{\delta} + \frac{12C_4}{\delta^2} \left(\frac{y}{\delta}\right)^2$$

Applying (iv) boundary condition i.e., $\frac{\partial^2 u}{\partial y^2} = 0$ at $y = 0$, we obtain

$$0 = C_2 \quad (3)$$

Applying (v) boundary condition i.e., $\frac{\partial^2 u}{\partial y^2} = 0$ at $y = \delta$, we obtain

$$0 = C_3 + 2C_4 \quad (4)$$

Solving Eqs (1), (2), (3) and (4) simultaneously, we obtain

$$C_1 = 2, \quad C_3 = -2 \quad \text{and} \quad C_4 = 1$$

Then the velocity profile becomes

$$\frac{u}{U_\infty} = 2 \frac{y}{\delta} - 2 \left(\frac{y}{\delta}\right)^3 + \left(\frac{y}{\delta}\right)^4$$

Q4.

The most general sinusoidal velocity profile for laminar boundary-layer flow on a flat plate is $u = A \sin(By) + C$. State three boundary conditions applicable to the laminar boundary-layer velocity profile. Evaluate constants A , B and C .

For the above velocity profile find expressions for:

- the rate of growth of δ as a function of x .
- the local skin friction coefficient in terms of distance and flow properties.
- the total drag force on a plate of length L and width w .

Solution

The velocity profile is given as

$$u = A \sin(By) + C$$

The boundary conditions are

- at $y = 0$, $u = 0$ (no-slip condition at the plate)
- at $y = \delta$, $u = U_\infty$ (free stream velocity at the edge of the boundary layer)
- at $y = \delta$, $\frac{\partial u}{\partial y} = 0$ (zero shear stress at the edge of the boundary layer)

Applying $u = 0$ at $y = 0$, one can get

$$0 = C$$

Applying $u = U_\infty$ at $y = \delta$ gives

$$U_\infty = A \sin(B\delta) \quad (1)$$

Applying (iii) boundary condition i.e., $\frac{\partial u}{\partial y} = 0$ at $y = \delta$, we get

$$0 = AB \cos(B\delta)$$

or
$$B\delta = \frac{\pi}{2}$$

or
$$B = \frac{\pi}{2\delta}$$

Substituting $B = \frac{\pi}{2\delta}$ in Eq. (1), we obtain

$$U_\infty = A \sin\left(\frac{\pi}{2\delta} \delta\right) = A \sin\left(\frac{\pi}{2}\right) = A$$

or
$$A = U_\infty$$

The velocity profile is then

$$u = U_\infty \sin\left(\frac{\pi y}{2\delta}\right)$$

Substituting $\frac{u}{U_\infty} = \sin\left(\frac{\pi y}{2\delta}\right)$ into von Karman integral equation

$\left(\frac{\tau_w}{\rho U_\infty^2} = \frac{d}{dx} \int_0^\delta \frac{u}{U_\infty} \left(1 - \frac{u}{U_\infty}\right) dy\right)$, we obtain

$$\begin{aligned} \frac{\tau_w}{\rho U_\infty^2} &= \frac{d}{dx} \int_0^\delta \frac{u}{U_\infty} \left(1 - \frac{u}{U_\infty}\right) dy \\ &= \frac{d}{dx} \int_0^\delta \sin\left(\frac{\pi y}{2\delta}\right) \left[1 - \sin\left(\frac{\pi y}{2\delta}\right)\right] dy \\ &= \frac{d}{dx} \int_0^\delta \left[\sin\left(\frac{\pi y}{2\delta}\right) - \sin^2\left(\frac{\pi y}{2\delta}\right)\right] dy \\ &= \frac{d}{dx} \int_0^\delta \left[\sin\left(\frac{\pi y}{2\delta}\right) - \frac{1}{2} \left\{1 - \cos\left(\frac{\pi y}{\delta}\right)\right\}\right] dy \quad \left[\because \sin^2 \theta = \frac{1 - \cos 2\theta}{2}\right] \\ &= \frac{d}{dx} \int_0^\delta \left[\sin\left(\frac{\pi y}{2\delta}\right) - \frac{1}{2} + \frac{1}{2} \cos\left(\frac{\pi y}{\delta}\right)\right] dy \\ &= \frac{d}{dx} \left[\frac{-\cos\left(\frac{\pi y}{2\delta}\right)}{\frac{\pi}{2\delta}} - \frac{1}{2} y + \frac{1}{2} \frac{\sin\left(\frac{\pi y}{\delta}\right)}{\frac{\pi}{\delta}} \right]_0^\delta \end{aligned}$$

$$\begin{aligned}
&= \frac{d}{dx} \left[\frac{-\cos\left(\frac{\pi \delta}{2\delta}\right)}{\frac{\pi}{2\delta}} + \frac{\cos\left(\frac{\pi \times 0}{2\delta}\right)}{\frac{\pi}{2\delta}} - \frac{1}{2}\delta - \frac{1}{2} \times 0 + \frac{\sin\left(\frac{\pi \delta}{\delta}\right)}{\frac{2\pi}{\delta}} - \frac{\sin\left(\frac{\pi \times 0}{\delta}\right)}{\frac{2\pi}{\delta}} \right] \\
&= \frac{d}{dx} \left[0 + \frac{1}{\frac{\pi}{2\delta}} - \frac{1}{2}\delta + 0 - 0 \right] \\
&= \frac{d}{dx} \left[\frac{2\delta}{\pi} - \frac{\delta}{2} \right] \\
&= \frac{d\delta}{dx} \left[\frac{2}{\pi} - \frac{1}{2} \right] = 0.137 \frac{d\delta}{dx}
\end{aligned}$$

or
$$\tau_w = 0.137\rho U_\infty^2 \frac{d\delta}{dx} \quad (2)$$

The velocity at a location y measured from the plate is $u = U_\infty \sin\left(\frac{\pi y}{2\delta}\right)$

Differentiating with respect to y , we have

$$\frac{\partial u}{\partial y} = U_\infty \frac{\pi}{2\delta} \cos\left(\frac{\pi y}{2\delta}\right)$$

At $y = 0$,

$$\left. \frac{\partial u}{\partial y} \right|_{y=0} = U_\infty \frac{\pi}{2\delta} \cos\left(\frac{\pi \times 0}{2\delta}\right) = \frac{\pi U_\infty}{2\delta}$$

Therefore, the wall shear stress becomes

$$\tau_w = \mu \left. \frac{\partial u}{\partial y} \right|_{y=0} = \mu \frac{\pi U_\infty}{2\delta} = \frac{\mu \pi U_\infty}{2\delta} \quad (3)$$

Equating the values of τ_w from Eqs (2) and (3), we obtain

$$\tau_w = 0.137\rho U_\infty^2 \frac{d\delta}{dx} = \frac{\mu \pi U_\infty}{2\delta}$$

or,

$$\delta \frac{d\delta}{dx} = \frac{\pi \mu}{0.274\rho U_\infty}$$

or,

$$\delta d\delta = 11.466 \frac{\mu}{\rho U_\infty} dx$$

Integrating the above equation, we get

$$\frac{\delta^2}{2} = 11.466 \frac{\mu}{\rho U_\infty} x + C$$

At $x = 0$, $\delta = 0$, so $C = 0$.

Then

$$\frac{\delta^2}{2} = 11.466 \frac{\mu}{\rho U_\infty} x$$

or,
$$\delta^2 = \frac{22.932\mu}{\rho U_\infty} x$$

or,
$$\delta = \sqrt{\frac{22.932\mu}{\rho U_\infty} x} = 4.79 \sqrt{\frac{\mu}{\rho U_\infty} x}$$

or,
$$\frac{\delta}{x} = 4.79 \sqrt{\frac{\mu}{\rho U_\infty x}} = \frac{4.79}{\sqrt{\text{Re}_x}}$$

The local *skin friction coefficient* C_{fx} is given by

$$\begin{aligned} C_{fx} &= \frac{\tau_w}{\frac{1}{2}\rho U_\infty^2} \\ &= \frac{\mu\pi U_\infty}{2\delta} = \frac{\mu\pi}{\rho\delta U_\infty} = \frac{\mu\pi}{\rho U_\infty \times 4.79 \sqrt{\frac{\mu}{\rho U_\infty} x}} \\ &= \frac{0.656}{\sqrt{\text{Re}_x}} \end{aligned}$$

The total shear force F_D on one side of the plate is given by

$$\begin{aligned} F_D &= \int_0^L \tau_w w dx \\ &= \int_0^L \frac{\pi\mu U_\infty}{2\delta} w dx \\ &= \int_0^L \frac{\pi\mu U_\infty}{2 \times 4.79 \sqrt{\frac{\mu}{\rho U_\infty} x}} w dx \\ &= 0.327 \rho U_\infty^2 w \sqrt{\frac{\mu}{\rho U_\infty}} \int_0^L \frac{1}{\sqrt{x}} dx \\ &= 0.327 \rho U_\infty^2 w \sqrt{\frac{\mu}{\rho U_\infty}} \left[2x^{\frac{1}{2}} \right]_0^L \\ &= 0.655 \rho U_\infty^2 w \sqrt{\frac{\mu}{\rho U_\infty}} L^{\frac{1}{2}} \end{aligned}$$